

Computing the Poisson kernel through Fourier transform (in the "I know the trick manner")

We seek $G(z)$. A.T. $\Delta G(z) = \delta^3(z)$ (*)

Set $G(z) = \int_{\mathbb{R}^3} \hat{G}(k) e^{-i k \cdot z} d^3 k$

$\Leftrightarrow \hat{G}(k) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} G(z) e^{i k \cdot z} d^3 z$

$\delta^3(z) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{-i k \cdot z} d^3 k$

plugging that in (*) yields

$\Delta \int_{\mathbb{R}^3} \hat{G}(k) e^{-i k \cdot z} d^3 k = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{-i k \cdot z} d^3 k$

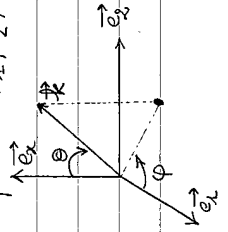
$\int_{\mathbb{R}^3} (-i)^2 \|\hat{k}\|^2 \hat{G}(k) e^{-i k \cdot z} d^3 k = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{-i k \cdot z} d^3 k$

stripping the integral and equating (that is up to null measure modifications) yields:

$\hat{G}(k) = \frac{-1}{(2\pi)^3} \times \frac{1}{\|\hat{k}\|^2} \text{ a.e.}$

More importantly $G(z) = \frac{-1}{(2\pi)^3} \int_{\mathbb{R}^3} \frac{1}{\|\hat{k}\|^2} e^{-i k \cdot z} d^3 k$.

Let us now choose a system of coordinates with respect to the frame $(\vec{e}_1, \vec{e}_2, \vec{e}_3)$



Using spherical coordinates yields, one gets $k = \|k\|$,

$$G(z) = \frac{-1}{(2\pi)^3} \int_0^{+\infty} \int_0^\pi \int_0^{2\pi} \frac{1}{k^2} e^{-i k \|z\| \cos \theta} k^2 \sin \theta d\varphi d\theta dk$$

$$= \frac{-1}{(2\pi)^2} \int_0^{+\infty} \int_0^\pi e^{-i k \|z\| \cos \theta} \sin \theta d\theta dk$$

Let say, $u = -\cos \theta$ $u|_{\theta=0} = -1$
 $du = \sin \theta d\theta$ $u|_{\theta=\pi} = 1$

$\Rightarrow G(z) = \frac{-1}{(2\pi)^2} \int_0^{+\infty} \int_{-1}^1 e^{+i k \|z\| u} du dk$

$= \frac{-1}{(2\pi)^2} \int_0^{+\infty} \frac{1}{\|z\|} (e^{i k \|z\|} - e^{-i k \|z\|}) dk$

$= \frac{i}{(2\pi)^2} \times \frac{1}{\|z\|} \int_0^{+\infty} \frac{1}{k} (e^{i k \|z\|} - e^{-i k \|z\|}) dk$

Note that $\int_0^{+\infty} \frac{1}{k} e^{-i k \|z\|} dk = \int_0^\infty \frac{-1}{k} e^{+i k \|z\|} (-dk) = -\int_0^\infty \frac{1}{k} e^{i k \|z\|} dk$

Meaning that $G(z) = \frac{i}{(2\pi)^2} \frac{1}{\|z\|} \int_{\mathbb{R}} \frac{1}{k} e^{i k \|z\|} dk$

It is left to do is to compute the integral

$I = \int_{\mathbb{R}} \frac{1}{k} e^{i k x} dk$ $x = \|z\| > 0$

A solution is to write $I = \lim_{\epsilon \rightarrow 0^+} I_\epsilon = \lim_{\epsilon \rightarrow 0^+} \left(\int_{-\infty}^{-\epsilon} + \int_{\epsilon}^{+\infty} \right) \frac{1}{k} e^{i k x} dk$

and one uses a complex analysis trick.

Consider $k \in \mathbb{C}$, note that $e^{i k x} = e^{i \text{Re}(k)x} e^{-\text{Im}(k)x}$

and consider the contour , no pole is enclosed

by the contour \Rightarrow the overall integral is null and one

gets $0 = \int_{\Gamma} \dots \rightarrow \dots + \underbrace{\int_{\Gamma} \dots}_{\text{due to exponential decay}} + \underbrace{\int_{\Gamma} \dots}_{\text{decay}}$

The integral on Γ_R vanishes and one gets

$$I_{\epsilon} = \dots = -\underbrace{\int_{\Gamma} \dots}_{\epsilon} + \underbrace{\int_{\Gamma} \dots}_{\epsilon}$$

meaning $I = \lim_{\epsilon \rightarrow 0^+} \int_{\Gamma} \frac{1}{k} e^{i\|\vec{x}\|} dk$, on Γ_{ϵ} one has $k = \epsilon e^{i\theta}$, $0 \leq \theta \leq \pi$

$$= \lim_{\epsilon \rightarrow 0^+} \int_0^{\pi} \frac{1}{\epsilon e^{i\theta}} e^{i\epsilon e^{i\theta} \|\vec{x}\|} i \epsilon e^{i\theta} d\theta$$

$$= \lim_{\epsilon \rightarrow 0^+} i \int_0^{\pi} e^{i\epsilon e^{i\theta} \|\vec{x}\|} d\theta$$

$= i\pi$, thanks to the limit which kills higher order terms.

Finally one gets $G(\vec{x}) = \frac{-1}{4\pi} \times \frac{1}{\|\vec{x}\|}$

up to solutions of homogeneous equation.

Remark that the integration over the contour Γ yields the same result for $I = \int_{\mathbb{R}} \frac{1}{k} e^{ikx} dk$, this is not obvious, Green's functions are known to have various determinations (retarded/advanced solutions Any).